



UNIVERSITÄT  
HEIDELBERG  
ZUKUNFT  
SEIT 1386

# Unstable domain walls and transitions in the flux Landscape

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String Phenomenology, Boston, July 8, 2025

Work in progress with B. Friedrich, A. Hebecker and M. Wiesner

# Flux Landscape and domain walls

→ Landscape approach to naturalness issues [Bousso, Polchinski '00][Denef, Douglas '04] [...]

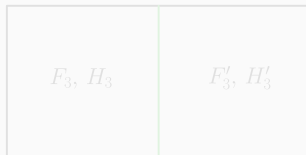
→ IIB dS solutions KKLT and LVS [Kachru, Kallosh, Linde, Trivedi '03][Balasubramanian, Berglund, Conlon, Quevedo '05] ...

Concerns: [almost all String Pheno attendees '18–Now]

But in this talk: we take the dS Landscape for granted

→ Focus on transitions and domain walls: stability?

D5/NS5-branes on three-cycles of CY



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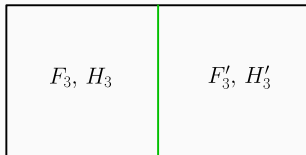
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# Outline

1/ Domain walls are **unstable**: no instanton

[Johnson, Larfors '08]×2 [Aguirre, Johnson, Larfors '10]

→ Field theory toy model

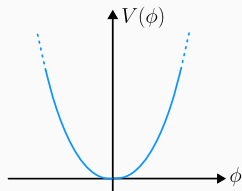
→ KKLT and LVS

→ Gravity

2/ Other ways to transition and how to estimate the rate?

# Field theory toy model

Scalar potential:  $V(\phi) = V_0 + \frac{m^2}{2}\phi^2 + \mathcal{O}(\phi^3)$



Brane potential:  $V_{\text{brane}}(\phi) = Ae^{-B\phi}$ ,  $B = \mathcal{O}(1)$

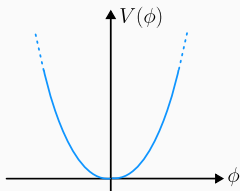
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→ Same vacuum both sides of DW ( $x = 0$ )       $\implies$  profile?

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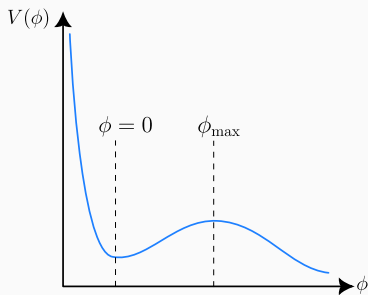
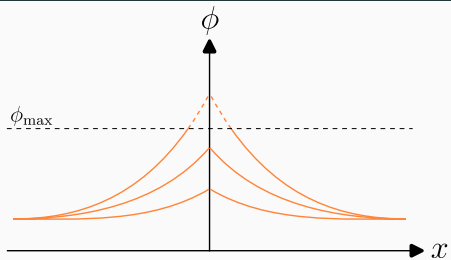
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# Intuition



# Profile

$$\phi'' = \frac{d}{d\phi}(V + V_{\text{brane}}) = m^2\phi - AB\delta(x)e^{-B\phi}$$

Derivative jump:  $\phi'(x)|_{-\epsilon}^{+\epsilon} = -ABe^{-B\phi(0)}$

Boundary conditions:  $\phi'(0_-) = -\phi'(0_+)$  and  $\phi'(x \rightarrow \pm\infty) = 0$

$$\implies \phi(x) = \phi(0)e^{-m|x|}$$

$$\longrightarrow m\phi(0) = AB e^{-B\phi(0)}$$

Large displacement if  $AB^2/m \gg 1$  and then  $\phi(0) \simeq \frac{1}{B} \ln\left(\frac{AB^2}{m}\right)$   
 $\simeq \ln\left(\frac{A}{m}\right)$

$\longrightarrow$  Do we have  $\phi_{\text{max}} > \ln(A/m)$  in KKLT and LVS?

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- KKLT

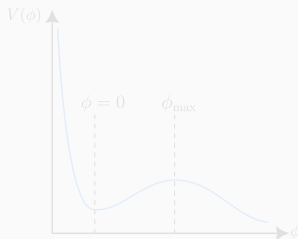
$$V_{\text{KKLT}} = \frac{aCg_s e^{-a\sigma}}{2\sigma^2} \left( \frac{a\sigma}{3} C e^{-a\sigma} - |W_0| + C e^{-a\sigma} \right) \quad \sigma = e\sqrt{\frac{2}{3}}\varphi$$

$$W_0 = -C e^{-a\sigma_0} \left( 1 + \frac{2}{3} a\sigma_0 \right) \ll 1 \quad \text{DW tension: } T_{\text{brane}} = \frac{A_0}{\mathcal{V}}$$

$$m_{\text{KKLT}} \mathcal{V}_0 \sim g_s^{\frac{1}{2}} |W_0| \frac{(-\ln |W_0|)}{2} \ll 1 \quad \text{regime OK!}$$

$$\rightarrow \text{Displacement: } \phi^{\text{KKLT}}(0) \simeq -\frac{2}{\sqrt{6}} \ln |W_0|$$

To be compared with  $\phi_{\text{max}}^{\text{KKLT}}$



- KKLT

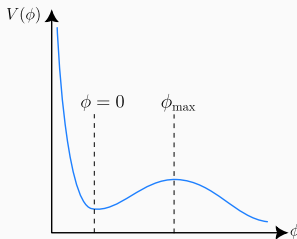
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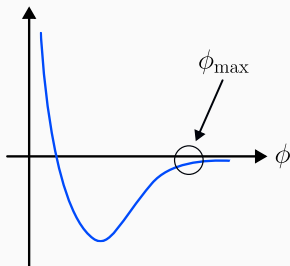
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→ Barrier width: from pre-uplift pot. [Westphal '08]



Approx. vanishing at  $\sigma = \sigma_0 + 1/a \implies \phi_{\max}^{\text{KKLT}} \simeq \sqrt{\frac{3}{2}} \frac{1}{(-\ln|W_0|)}$

Domain wall parametrically unstable

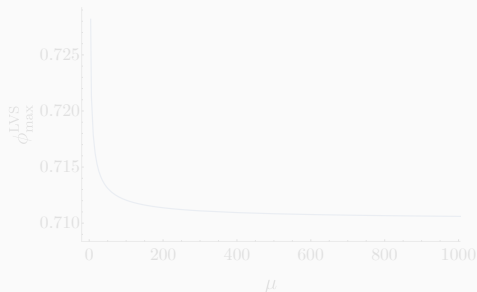
$$V(\mathcal{V}) \simeq \frac{g_s}{\mathcal{V}^3} \left[ \mu - \ln^{3/2}(\mathcal{V}) \right]$$

$$\implies m_{\text{LVS}} \mathcal{V}_0 \simeq \frac{g_s^{1/2}}{\mathcal{V}_0^{1/2}} \ll 1 \quad \text{regime OK!}$$

→ Displacement:  $\phi^{\text{LVS}}(0) \simeq -\frac{1}{\sqrt{6}} \ln \mathcal{V}_0$

→ Barrier width: Constant in large  $\mu$  limit

For a  $\overline{\text{D3}}$  uplift  $\implies \phi_{\text{max}}^{\text{LVS}} \approx 0.71$



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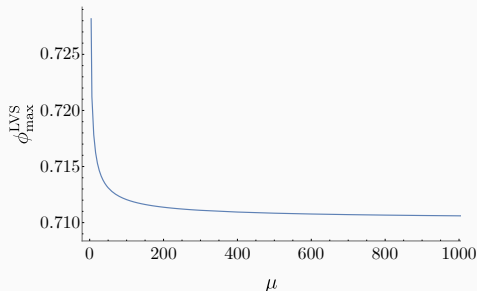
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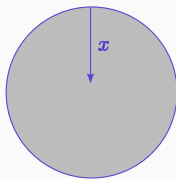
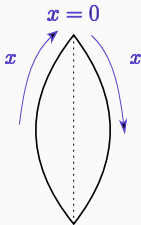
Domain wall  
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## Including gravity

$O(4)$  symmetric ansatz  $ds_4^2 = dx^2 + \rho(x)^2 d\Omega_3^2$  [Coleman, De Luccia '80]

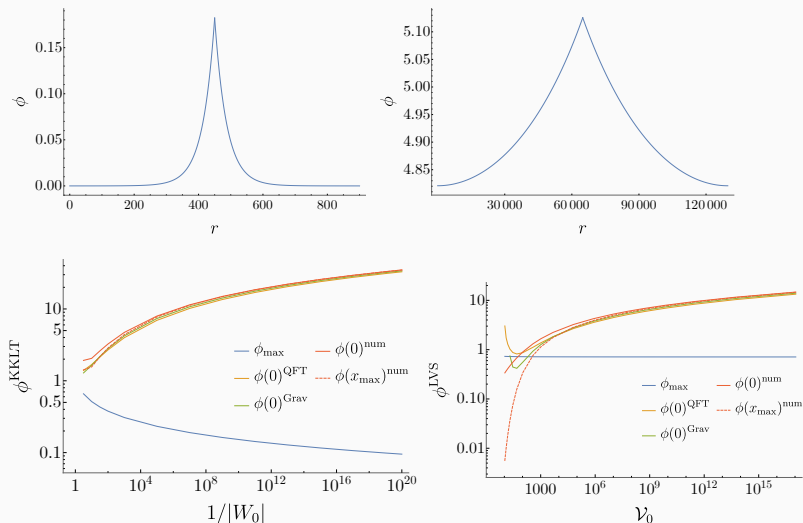
$$\rho' = \pm \sqrt{1 + \frac{\rho^2}{3} \left( \frac{\phi'^2}{2} - V \right)}, \quad \rho'' = -\frac{\rho}{3} \left( \phi'^2 + V + \delta(x) \frac{3}{2} V_{\text{brane}} \right),$$

$$\phi'' = m^2 \phi - 3 \frac{\rho'}{\rho} \phi' - \delta(x) A B e^{-B\phi},$$



Behaviour similar to field theory:  $\phi(0) \gtrsim \frac{1}{2B} \ln \left( \frac{3A^2 B^2}{4m^2} \right)$

Membrane DW profiles  $\rightarrow$  Undershoot/overshoot algorithm



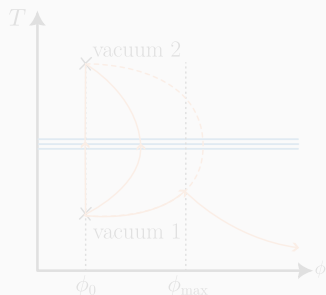
# Non-tunneling

- Studied in [Brown, Dahlen '11], flyover type: [Blanco-Pillado, Deng, Vilenkin, '19]

→ In presence of membrane?

- Tachyon for brane nucleation: [Sen '98][Shiu, Tonioni, Van Hemelryck, Van Riet '23 '24]

$$V(\phi, T) = V(\phi) + Ae^{-B\phi} \frac{1}{\cosh T}$$



Flyover channel?

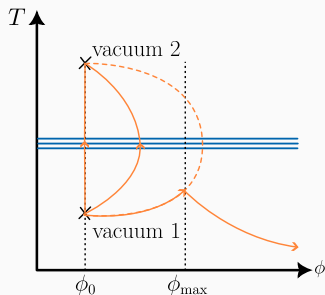
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# High tension causes problems

- dS with potential  $V_h \implies$  horizon  $l_h^{-1} \sim V_h^{1/2} M_P$

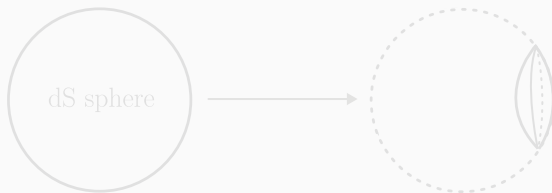
**Crit. radius**  $\rho_b \sim T_{\text{brane}}/\epsilon$  with  $T_{\text{brane}} \sim \mathcal{V}_0^{-1} M_P^3$  and  $\epsilon \sim V_h M_P^4$

Asking  $\rho_b < l_h \implies \frac{1}{\mathcal{V}_0} < V_h^{1/2}$

But in KKLT and LVS:  $V_h^{1/2} \sim m$  and  $m\mathcal{V}_0 \ll 1$

- Bubble mass  $T_{\text{brane}}\rho_b^2$

Compare with Nariai BH mass  $l_h M_P^2$ :  $\frac{T_{\text{brane}}\rho_b^2}{l_h M_P^2} \sim \frac{1}{(\mathcal{V}_0 m)^3} \gg 1$



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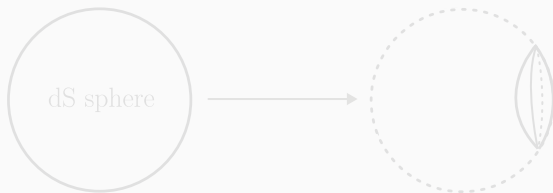
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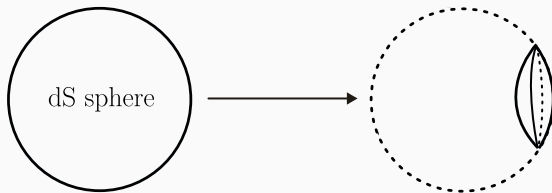
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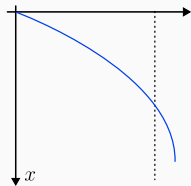
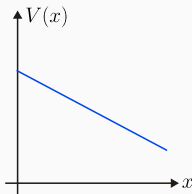
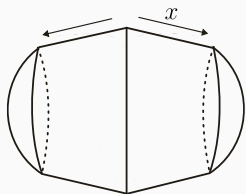
# Lorentzian approach

Wheeler-DeWitt equation [deWitt '67]

Hamiltonian formalism [Dirac '58][ADM '62][Fischler, Morgan, Polchinski '90]...[Céspedes, de Alwis, Muia, Quevedo '20]

$$ds^2 = -(N^t dt)^2 + L^2(dr^2 + N^r)^2 + R^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

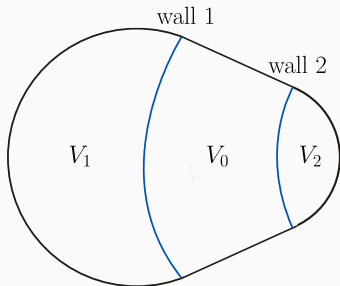
- Suppose we nucleate the bubble



dS patch is safe if it can reach horizon size!

# Two-wall system

[Kolitch, Eardley '97]



$$S = \int dt \left[ p_1 \dot{r}_1 + p_2 \dot{r}_2 + \int dr \left( \pi_L \dot{L} + \pi_R \dot{R} - N^t \mathcal{H}_t - N^r \mathcal{H}_r \right) \right]$$

→ Fix **radial param.** and impose **slicing condition**

→ **Solve constraints**  $\mathcal{H}_t = \mathcal{H}_r = 0$  in  $V_1, V_0, V_2$

# Conclusions

- The would-be ST dS Landscape **does not admit CdL transitions** [Aguirre, Johnson, Larfors '10]
- KKLT and LVS barrier width **too small** compared to field displacement at the DW
  - Asympt. regime in field theory, + gravity
  - Numerics for fully general case
- Apply to all “two-step constructions” (CS + Kähler)  $\implies$  F-term uplift?
- How to populate anyway?
  - Usual non-tunneling fail
  - Hamiltonian approach and WDW for probability of the dS patch to reach horizon size, work in progress, stay tuned!

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